### A New Conjugate Gradient Method and Application to Dynamic Load Identification Problems

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In this paper, a modified conjugate gradient (MCG) algorithm is proposed for solving the force reconstruction problems in practical engineering. This new method is derived from a stable regularization operator and is also strictly proved using the mathematical theory. Moreover, we also prove the sufficient descent and global convergence characteristic of the newly developed algorithm. Finally, the proposed algorithm is applied to force reconstruction for the airfoil structure and composite laminated cylindrical shell. Numerical simulations show that the proposed method is highly efficient and has robust convergence performances. Additionally, the accuracy of the proposed algorithm in identifying the expected loads is satisfactory and acceptable in practical engineering.

#### 1. INTRODUCTION

In recent years, force identification in structural dynamics has attracted lots of interest of many scholars.<sup>1-7</sup> The inverse force identification by structural vibration data such as displacement responses is particularly suitable for structural mechanics and structural health monitoring in the field of mechanical engineering. However, for a great many practical engineering problems, we have to try our best to know the corresponding information of external loads. For example, after knowing external force acting on the structure of aircraft wings and wind turbines, it is possible or convenient to exploit a great many advanced computational algorithms to ensure their safety. Also, we often optimize these structures to satisfy the requirements of industrial development. Moreover, in many cases, it is not easy to get the accurately applied loads thatact on engineering structures, so efforts are made to develop indirect computational inverse technology which computes the excitation loads with known vibration data, such as displacement responses, strain data, and acceleration and velocity responses.<sup>8–10</sup>

Jacquelin et al. analyzed a deconvolution technique and successfully used this method to recover an experimental force. <sup>11</sup> The finite element method is exploited to reconstruct the moving loads acting on the bridge deck. <sup>12</sup> Gunawan et al. proposed the Two-step B-splines regularization method and used

it to regularize the ill-posed problem of impact load reconstruction. 13 Zhu and Law exploited a new algorithm that is based on the regularization method and modal superposition and applied it for the moving force reconstruction.<sup>14</sup> A regularization method using signal processing techniques was proposed to deal with force identification.<sup>15</sup> The dynamic programming algorithm was exploited by Law et al. to identify the moving force. 16 Zhang and Ohsaki transformed the force reconstruction for the prestressed pin-jointed structure into an optimization problem and solved it using the method of simulated annealing method. 17 Xie et al. proposed an identification method based on statistical energy analysis to solve the high-frequency load identification.<sup>18</sup> Li et al. presented a new method that uses wavelet multi-resolution analysis to solve the load identification.<sup>19</sup> A novel load identification method was proposed to deal with the identification of discontinuous loading.<sup>20</sup>

However, there are still some shortcomings that exist in the research study so far. Firstly, the results of most traditional methods are weak anti-noise in identifying multi-source dynamic loads and are still not satisfactory when the noisy level increases. Additionally, we often use the iterative regularization method to obtain the solution of large-scale inverse problems, but their convergence rate in obtaining a regularized solution is pretty slow and inefficient.<sup>21</sup> Furthermore, very few

references are found about the load identification in the aspect of large noisy levels.<sup>22–24</sup> Therefore, we have to seek new advanced computational methods to overcome these shortcomings. A new fast and efficient conjugate gradient method is presented in this paper, also proved by mathematical theory, and exploited to reconstruct dynamic load in the time domain from measured displacement responses.

This paper proceeds as follows. Section 2 devotes to the establishment of the new fast and efficient conjugate gradient algorithm. The sufficient descent characteristic and global attractivity of the proposed method are strictly proved in Section 3. In Section 4, we present numerical studies to investigate the stability, accuracy, and effectiveness of the newly developed method. The performances of the proposed method under different measurement noisy levels are investigated in detail. Finally, the last section presents the conclusions of this paper.

# 2. THE CONSTRUCTION OF A NEW FAST AND EFFICIENT CONJUGATE GRADIENT ALGORITHM

Firstly, the following unconstrained optimization problem is investigated:

$$\min_{x \in P_n} f(x); \tag{1}$$

in which  $f: \mathbb{R}^n \to \mathbb{R}$  is a continuous and differentiable function. At point  $x_k$ , its gradient is denoted by  $g(x_k)$ . The corresponding iteration form is given as

$$x_{k+1} = d_k \alpha_k + x_k, k = 0, 1, 2 \cdots;$$
 (2)

in which  $\alpha_k$  denotes the steplength and the search direction is denoted by  $d_k$ . Usually, the one-dimensional search method is used to search the stepsize, and the corresponding formula is defined as

$$f(d_k\alpha_k + x_k) = \min_{\alpha \ge 0} f(d_k\alpha_k + x_k); \tag{3}$$

where  $d_k$  represents the search direction given by the following equation:

$$d_k = \begin{cases} -g_0, & k=0, \\ d_{k-1}\beta_k - g_k, & k \ge 1 \end{cases}; \tag{4}$$

in which  $\beta_k$  is a scalar, and scientific and reasonable choices for this parameter correspond to different kinds of new conjugate gradient methods. Herein, a new formula for this parameter is given by

$$\beta_k^{WDX} = \begin{cases} \frac{g_k^T d_k}{g_k^\alpha d_{k-1}} (\alpha \ge 1) & \text{if } ||g_k|| \ge 1\\ 1 & \text{otherwise} \end{cases} . \tag{5}$$

The corresponding new algorithm is given as the following steps:

#### MCG Algorithm

Step 0: Considering  $x_0 \in \mathbb{R}^n$ , set k = 0.

Step 1: Obtain  $\beta_k$  using the formula Eq. (5).

Step 2: Compute  $d_k$  on the basis of Eq. (4). When  $\|g_k\| < \varepsilon$ , then stop.

Step 3: Generate  $\alpha_k$  exploiting Eq. (3).

Step 4: Renewing next point according to Eq. (2). We will stop it when  $||g_k|| < \varepsilon$  and  $f(x_k) > f(x_{k+1})$  or else go to Step 0 with k = k + 1.

#### 3. GLOBAL CONVERGENCE

The validity and convergence of new parameter  $\beta_k^{WDX}$  will be investigated in the following part. To prove the good definition of the corresponding new algorithm for the new parameter, we will investigate its global convergence.

**Theorem 3.1.** Considering the conjugate gradient algorithm based on the new parameter  $\beta_k^{WDX}$ , Eqs. (3) and (4), then we have that there is C>0 such that

$$g_k^T d_k \le -C \|g_k\|^2; \tag{6}$$

when  $k \geq 0$ .

**Proof.** Firstly, we can easily obtain the assertion when k = 0. Secondly, the statement that the sufficient condition is also true for  $k \ge 1$  will be strictly proved.

Using Eq. (4), we can get

$$g_{k+1}^{T}d_{k+1} = g_{k+1}^{T}(-g_{k+1}+\beta_{k+1}d_{k}) = -\|g_{k+1}\|^{2} + \beta_{k+1}g_{k+1}^{T}d_{k}.$$
 (7)

We also easily obtain that  $g_{k+1}^T d_k = 0$  according to the exact line search. Then, we can obtain that

$$g_{k+1}^T d_{k+1} = -\|g_{k+1}\|^2$$
.

Therefore, we can assert that  $d_{k+1}$  is a feasible direction. Thus, the proof of Theorem 3.1 is completed.

**(H1).** For the level set  $R^n$ , there is a lower bound for f; For  $x_0$  given, there exists a neighborhood N of  $\Gamma = \{x \in R^n | f(x_0) \ge f(x)\}$  in which f is continuously differentiable.

**(H2)**. There is L > 0 such that

$$||q(y) - q(x)|| \le L||y - x||, \forall x, y \in N;$$

i.e., g(x) is Lipschitz continuous.

Exploiting (H1) and (H2), we can get the following assertions [25, 26]:

**Lemma 3.1** Suppose (H1)-(H2) are satisfied. Considering any conjugate gradient method by Eq. (4), in which  $\alpha_k$  meets the precise minimization rule and  $d_k$  represents a descent search direction. Therefore,

$$\sum_{k=0}^{\infty} \frac{g_k^T d_k}{\|d_k\|^2} < \infty.$$

Based on Lemma 3.1, we get the following results.

**Theorem 3.2.** Let (H1)-(H2) and the descent condition are satisfied. Considering the conjugate gradient method given by Eqs. (2) and (4), in which  $\alpha_k$  is obtained based on Eq. (3). Therefore, we have that either

$$\lim_{k\to\infty} \|g_k\| = 0;$$

or

$$\sum_{k=0}^{\infty} \frac{\left(g_k^T d_k\right)^2}{\left\|d_k\right\|^2} < \infty.$$

**Proof.** Actually, the conclusion can be proved using the contradiction method. Therefore there is a positivie parameter  $\varepsilon$  which satisfies

$$\epsilon \le \|g_k\|. \tag{8}$$

According to Eq. (4), we can get

$$\beta_{k+1}d_k = d_{k+1} + g_{k+1}.$$

So.

$$\beta_{k+1}^{2} \|d_{k}\|^{2} - \|g_{k+1}\|^{2} - 2g_{k+1}^{T} d_{k+1} = \|d_{k+1}\|^{2}.$$
 (9)

Therefore.

$$\frac{\|d_{k+1}\|^{2}}{(g_{k+1}^{T}d_{k+1})^{2}} = \frac{\beta_{k+1}^{2}\|d_{k}\|^{2}}{(g_{k+1}^{T}d_{k+1})^{2}} - \left(\frac{1}{\|g_{k+1}\|} + \frac{\|g_{k+1}\|^{2}}{g_{k+1}^{T}d_{k+1}}\right)^{2} + \frac{1}{\|g_{k+1}\|^{2}} \le \frac{\beta_{k+1}^{2}\|d_{k}\|^{2}}{(g_{k+1}^{T}d_{k+1})^{2}} + \frac{1}{\|g_{k+1}\|^{2}} = \frac{2}{\|g_{k+1}\|^{2}}.$$
(10)

Then, we obtain

$$\sum_{i=0}^{K} \frac{2}{\|g_i\|^2} \ge \frac{2}{\|g_k\|^2} \ge \frac{\|d_k\|^2}{(g_k^T d_k)^2}.$$
 (11)

So,

$$\frac{\varepsilon^2}{2K} \le \frac{(d_k g_k)^2}{\|d_k\|^2}.\tag{12}$$

By using Eqs. (8) and (12), we have

$$\sum_{k=0}^{\infty} \frac{g_k^T d_k}{\|d_k\|^2} = \infty;$$

which immediately contradicts the conclusion of Lemma 3.1. Therefore, Theorem 3.2 is fully proved.

#### 4. APPLICATION

To validate the effectiveness and stability of MCG in reconstructing multi-source loads acting on an airfoil structure, the force reconstruction problem is investigated in the following subsection. Noticing the linear and time-invariant assumptions with regard to load identification, the corresponding convolution equation can be expressed as:<sup>27</sup>

$$y(t) = \int_0^t p(\tau)G(t-\tau)d\tau; \tag{13}$$

in which y(t) denotes the displacement response, G(t) represents the kernel of the impulse response, and p(t) represents the external dynamic force.

The corresponding equally spaced intervals can be obtained using the discretization of the time period, and then we give Eq. (13) in matrix form as: $^{28-30}$ 

$$G(t)P(t) = Y(t); (14)$$

or given by

$$\begin{pmatrix} y_1 \\ y_2 \\ \vdots \\ y_m \end{pmatrix} = \begin{pmatrix} g_1 \\ g_2 & g_1 \\ \vdots & \vdots & \ddots \\ g_m & g_{m-1} & \cdots & g_1 \end{pmatrix} \begin{pmatrix} p_1 \\ p_2 \\ \vdots \\ p_m \end{pmatrix} \triangle t.$$

Considering the static characteristic of the structure before the force is applied, we can assert that  $y_0 = 0$ ,  $g_0 = 0$ , and then G is the lower triangular matrix.

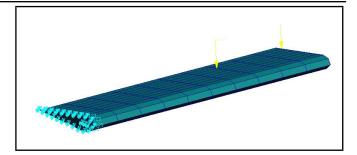


Figure 1. The finite element model of the airfoil.

In order to reconstruct P(t), it is very necessary to get the value of G(t) and y(t). Exploiting the finite element method, we can compute the response at any point and the numerical Green's function. More importantly, it is not easy to get the dynamic force P(t) by a direct inverse operation. In the following part, the proposed method will be used to solve this kind of ill-posed inverse problem about force reconstruction of an airfoil structure and a composite laminated cylindrical shell.

#### 4.1. An Airfoil Structure

A practical engineering problem is investigated to determine vertical forces acting on an airfoil structure as shown in Fig. 1. The density, Poisson's ratio, and elastic modulus of the airfoil structure are  $\rho=8.3\times10^3{\rm kg/m^3},~\nu=0.3,$   $E=3.8\times10^{11}{\rm MPa},$  respectively. On the outer surface of the airfoil structure, there is a vertical concentrated force. The corresponding displacement responses can also be vertically obtained. In addition, one side of the airfoil structure is fixed. The arrows in Fig. 1 denote the corresponding points at which the dynamic loads are acting on the airfoil structure.

The corresponding dynamic forces are expressed as the following formulas:

$$F_1(t) = \begin{cases} q_1 \sin(\frac{2\pi t}{t_d}), & 0 \le t \le 2t_d \\ 0, & t < 0 \text{ and } t > 2t_d \end{cases}$$

$$F_2(t) = \begin{cases} 4q_2t/t_d, & 0 \le t \le t_d/4 \\ 2q_2 - 4q_2t/t_d, & t_d/4 < t \le 3t_d/4 \\ 4q_2t/t_d - 4q_2, & 3t_d/4 < t \le t_d \\ 0, & t > t_d \end{cases}$$

in which the parameters  $q_i(i=1,2),t_d$  are 1000N, 800N, and 0.004s, respectively. The time histories of these two forces are drawn by the software MATLAB in Figs. 2 and 3, respectively. Moreover, the numerical method is used to simulate the response of the test data. Figures 4 and 5 show the displacement responses at nodes 391 and 640, respectively. For the simulation of a practical noisy measurement, the noisy response can be expressed as

$$Y_{err} = rand(-1, 1) \cdot std(Y_{cal}) \cdot l_{noise} + Y_{cal},$$

where  $Y_{cal}$  and  $std(Y_{cal})$  represent the computer generated response and its standard deviation, respectively.  $l_{noise}$  and rand(-1,1) represent the noisy level and a random number between [-1,1].

To investigate the reverse computing ability of Landweber, the original conjugate gradient method (FRCG), and MCG, different noise levels such as 5%, 10%, 20% and 50% are considered. The other parameters are:  $\alpha = 6$  and  $\epsilon = 0.33$ . The

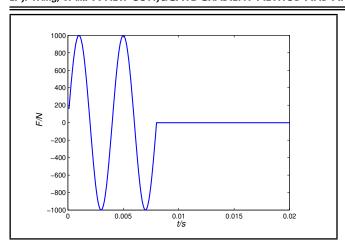


Figure 2. The vertical concentrated sine load acting on the outside surface.

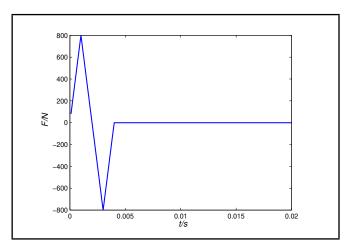


Figure 3. The vertical concentrated triangle load acting on the outside surface.

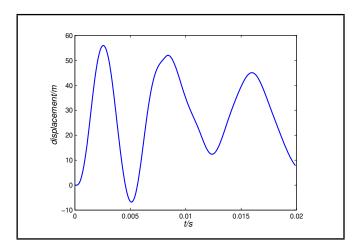


Figure 4. The corresponding vertical displacement response at node 391.

proposed algorithm is used to reconstruct the sine and triangle loads. The performances of the traditional Landweber iteration regularization method, FRCG, and MCG will be compared based on the following relative error

$$\tilde{F} = \|\frac{F_{Real} - F_{Identified}}{F_{Real}}\|.$$
 (15)

and

$$F_{\text{Average}} = \frac{1}{n} \sum_{i=1}^{n} \left| \frac{F_{\text{Real}}(i) - F_{\text{Identified}}(i)}{\max\{F_{j}\}} \right| * 100; \quad (16)$$

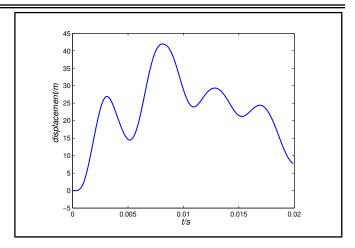


Figure 5. The corresponding vertical displacement response at node 640.

where  $i = 1, 2, \dots, n, j = 1, 2$ .

As shown in Fig. 6, under the condition of 5% noise level, Landweber, FRCG, and MCG make good performances in identifying the dynamic force. Landweber's and FRCG's performances are slightly worse than MCG's. Their relative deviations are displayed in Fig. 7. It can be found from these figures that Landweber's and FRCG's errors are much greater than MCG's, which shows the superiority of MCG. Furthermore, as can be seen from Table 1, this conclusion can also be drawn. This is mainly due to the advantage of MCG. As shown in Table 1, most deviations of Landweber and FRCG respectively mainly focus on the range of 11.58%, 9.80%, while most deviations of MCG concentrate in the range of 10.84%. Furthermore, Landweber's maximal deviation and average deviation in identifying the sine force are 7.46%, 0.91%, respectively. FRCG's maximal deviation and average deviation in identifying the sine force are 9.80%, 2.46%, respectively. The average deviation and maximal deviation of MCG in identifying the sine force are 0.88%, 5.22%, respectively. Additionally, this table also shows that Landweber's average and maximal deviation in triangular force recognition are 0.79\%, 11.58\%, respectively, and FRCG's average and maximal deviation in triangular force recognition are 1.58%, 8.30%, respectively. Those of MCG are 0.76%, 10.84%, respectively. At the same time, the iterative numbers of Landweber and FRCG are 49 and 26, respectively, while the iterative number of the proposed method is 23. In a word, these results above illustrate that the newly developed MCG method achieves good performance and gives stable, effective, and satisfactory results in reconstructing multisource dynamic loads. Therefore, at other different noise levels, we just compare the performances of the proposed method and the Landweber method in the next part.

As the noise level increases, the identified results slowly get worse, which can be shown in Figs. 8, 9, 10. Their relative deviations are respectively displayed in Figs. 11, 12, 13. It is also noted that the identified results under different noise levels are still good and also acceptable in engineering. These results also show that the proposed algorithm has a powerful identification ability.

As shown in Table 2, under 10% of noise level, most deviations of Landweber mainly focus on the range of 10.22%, while those of MCG concentrate in the range of 9.56%. Additionally, the maximum and mean error of Landweber are

**Table 1.** The identified force at five time points at noise level 5%.

			Landweber method		MCG method		FRCG method	
	Timepoint	Realforce	Identifiedforce	Error (%)	Identifiedforce	Error (%)	Identified force	Error (%)
Sine	0.001	1000	977.21	2.28	1003.1	0.31	1006.9	0.69
Triangle	0.0006	480	497.29	2.16	492.91	1.61	468.48	1.44
Sine	0.003	-1000	-982.62	1.74	-979.9	2.01	-1037.5	3.75
Triangle	0.001	800	707.37	11.58	713.25	10.84	733.58	8.30
Sine	0.0045	707.11	709.79	0.27	707.68	0.06	756.17	4.91
Triangle	0.0016	320	309.59	1.30	312.52	0.94	300.1	2.49
Sine	0.0063	-453.99	-460.98	0.70	-448.52	0.55	-445.51	0.85
Triangle	0.0033	-560	-568.11	1.01	-569.87	1.23	-565.14	0.64
Sine	0.0073	-891.01	-896.88	0.59	-877.55	1.35	-861.86	2.91
Triangle 0.0038 -160		-135.04	3.12	-134.83	3.15	-157.41	0.32	
Error (%)		Maximum	Average	Maximum	Average	Maximum	Average	
Sine		7.46	0.91	5.22	0.88	9.80	2.46	
Triangle		11.58	0.79	10.84	0.76	8.30	1.58	
Iterative steps		49		23		26		

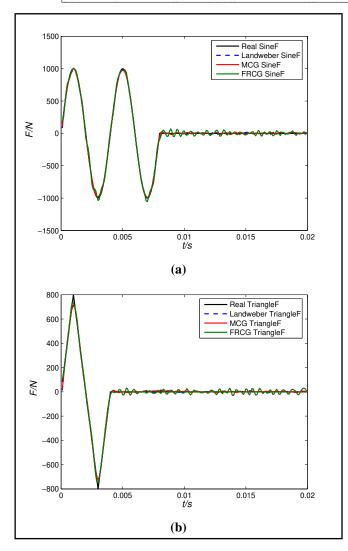
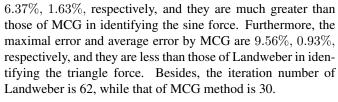
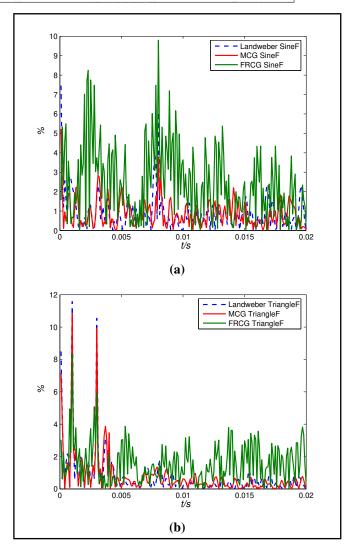


Figure 6. The identified sine and triangle force at noise level 5%; the number of iterations:  $N_{Landweber}=49, N_{MCG}=23.$ 



As shown in Table 3, under 20% of noise level, most deviations of MCG concentrate in the range of 13.16%, which is



**Figure 7.** The relative deviations for the identified sine and triangle force at noise level 5%.

smaller than Landweber. Additionally, the average and maximal error of Landweber in identifying the sine force are 3.78%, 13.74%, respectively, which are greater than those of MCG (3.32% and 13.16%, respectively). Moreover, in identifying the triangle force, it can also be shown that the average and maximal errors by Landweber are 2.16%, 8.98%, respectively, while those of the present method are 1.76%, 10.21%, respectively. However, the average error is less than that of Landwe-

**Table 2.** The identified force at five time points at noise level 10%.

	-			method	MCG method		
	Timepoint Realforce		Identifiedforce	Identifiedforce Error (%)		Error (%)	
Sine	0.001	1000	1024.8	2.48	997.68	0.23	
Triangle	0.0006	480	492.02	1.50	482.14	0.27	
Sine	0.003	-1000	-993.7	0.63	-1041.2	4.12	
Triangle	0.001	800	718.26	10.22	723.55	9.56	
Sine	0.0045	707.11	689.39	1.77	730.58	2.35	
Triangle	0.0016	320	330.82	1.35	305.01	1.87	
Sine	0.0063	-453.99	-453.61	0.04	-455.63	0.16	
Triangle	0.0033	-560	-587.92	3.49	-583.57	2.95	
Sine	0.0073	-891.01	-886.26	0.48	-897.66	0.67	
Triangle	Triangle 0.0038 -160		-151.99	1.00	-147.24	1.60	
Error (%)			Maximum	Average	Maximum	Average	
	Sine		6.37	1.63	5.13	1.46	
	Triangle		10.22	0.97	9.56	0.93	
	Iterative steps	S	62		30		

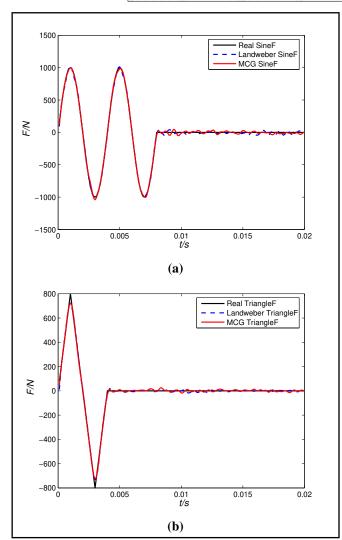


Figure 8. The identified sine and triangle force at noise level 10%; the number of iterations:  $N_{Landweber}=62, N_{MCG}=30.$ 

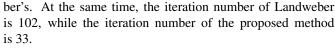


Table 4 displays the identified performances of Landweber and MCG in detail under 50% of the noise level. It shows that most errors of the present method are less than Landweber, which owes to the stable and efficient identification of the proposed algorithm. As shown in Table 4, most of Landwe-

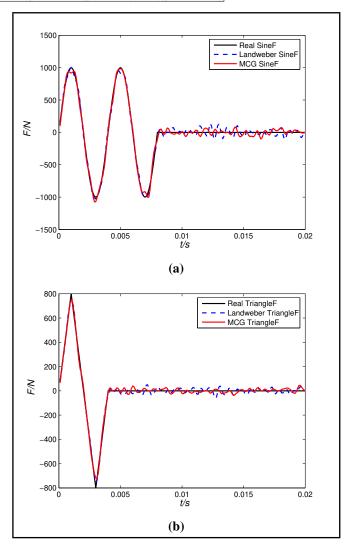
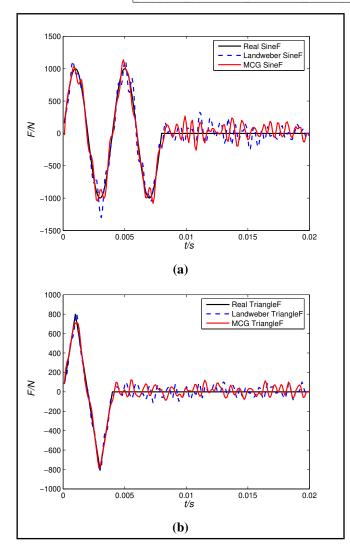


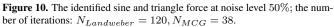
Figure 9. The identified sine and triangle force at noise level 20%; the number of iterations:  $N_{Landweber} = 102, N_{MCG} = 33.$ 

ber's deviations mainly focus on the range of 32.69%, while those of MCG concentrate in the range of 26.90%. Additionally, Landweber's maximum and mean deviations are 32.69%, 9.10%, respectively, much greater than those of MCG in identifying the sine force. Moreover, in the triangle force identification, the average and maximal deviation of MCG are 4.81%, and 15.27%, respectively, but MCG's average deviation is less than Landweber's. Besides, Landweber's iteration number

Table 3. The identified force at five time points at noise level 20%.

			Landweber	method	MCG method		
	Timepoint	Realforce	Identifiedforce	lentifiedforce Error (%)		Error (%)	
Sine	0.001	1000	1005.7	0.57	915.48	8.45	
Triangle	0.0006	480	476.16	0.48	468.07	1.49	
Sine	0.003	-1000	-1028.1	2.81	-1069.2	6.92	
Triangle	0.001	800	728.12	8.99	751.99	6.00	
Sine	0.0045	707.11	666.44	4.07	664.08	4.30	
Triangle	0.0016	320	282.81	4.65	273.89	5.76	
Sine	0.0063	-453.99	-463.72	0.97	-488.94	3.50	
Triangle	0.0033	-560	-591.35	3.92	-585.15	3.14	
Sine	0.0073	-891.01	-932.21	4.12	-1000.9	10.99	
Triangle	Triangle 0.0038 -160		-130.6	3.68	-145.77	1.78	
Error (%)			Maximum	Average	Maximum	Average	
	Sine		13.74	3.78	13.16	3.32	
	Triangle		8.98	2.16	10.21	1.76	
Iterative steps			102		33		





is 120, while that of the MCG method is 38. This further validates that the proposed method is much better.

## 4.2. A Composite Laminated Cylindrical Shell

Now we will consider a composite laminated cylindrical shell<sup>4</sup> as shown in Fig. 14. We will identify the impact forces acting on this structure by the proposed method. The mate-

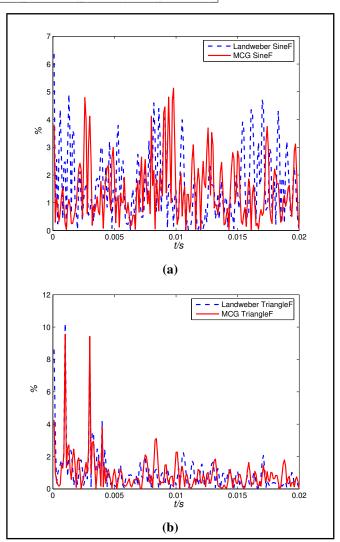
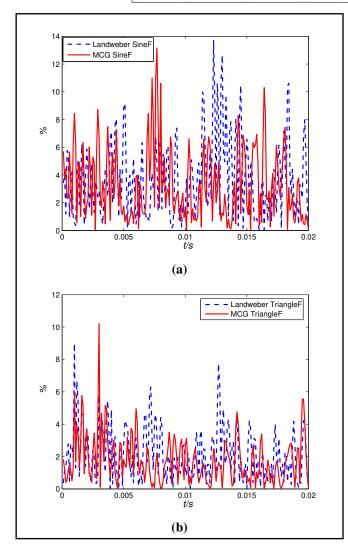


Figure 11. The identified sine and triangle force at noise level 50%; the number of iterations:  $N_{Landweber}=120, N_{MCG}=38.$ 

rial properties of the carbon/epoxy and glass/epoxy are shown in Table 5. The radial impact loads act on the outside surface. The measured displacement responses are along the radial direction. One end of the shell structure is free, and the other end is fixed. The arrow in Fig. 14 represents the action point of impact loads. Similarly, as in Section 4.1, the noisy responses with 10% of noise level are computed and displayed in Figs. 15, 16. We reversely compute the impact loads acting

**Table 4.** The identified force at five time points at noise level 50%.

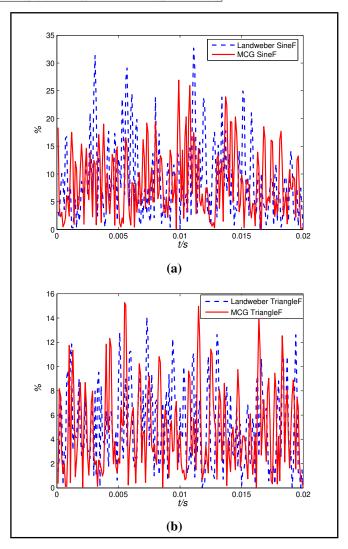
				method	MCG method		
	Timepoint Realforce		Identifiedforce	Error (%)	Identifiedforce	Error (%)	
Sine	0.001	1000	930.56	6.94	975.97	2.40	
Triangle	0.0006	480	503.65	2.96	463.91	2.01	
Sine	0.003	-1000	-1246.4	24.64	-875.46	12.45	
Triangle	0.001	800	759.42	5.07	706.05	11.74	
Sine	0.0045	707.11	778.07 7.10		571.92	13.52	
Triangle	0.0016	320	286.88	4.14	301.42	2.32	
Sine	0.0063	-453.99	-413.97	4.00	-449.66	0.43	
Triangle	0.0033	-560	-491.92	8.51	-561.1	0.14	
Sine	0.0073	-891.01	-863.44	2.76	-1082.8	19.18	
Triangle	Triangle 0.0038 -160		-130.02	3.75	-170.47	1.31	
Error (%)			Maximum	aximum Average Maximum		Average	
	Sine		32.69	9.10	26.90	8.47	
	Triangle		14.07	5.07 15.27		4.81	
Iterative steps			120	-	38		



**Figure 12.** The relative deviations for the identified sine and triangle force at noise level 20%.

on the composite laminated cylindrical shell structure by the Landweber, the original conjugate gradient method (FRCG), and the MCG, and investigate their reverse ability under a noise level of 10%. Additionally, their identified results will also be compared by Eqs. (15) and (16).

As shown in Figs. 17 and 18, under the condition of 10% of noise level, Landweber, FRCG, and MCG perform well in identifying the impact loads. Their relative deviations are dis-



**Figure 13.** The relative deviations for the identified sine and triangle force at noise level 50%.

played in Figs. 19 and 20. It can be seen from these figures that Landweber's and FRCG's error is much greater than MCG's, which shows the superiority of the MCG. As shown in Table 6, most deviations of Landweber and FRCG mainly focus on the range of 22.52%, 25.65%, while most deviations of the MCG concentrate in the range of 20.2%. Furthermore, Landweber's maximal deviation and average deviation in identifying the first impact force are 22.52%, 8.84%, respectively. FRCG's maximal

Table 5. The material properties of composite laminated cylindrical shell.

	Material properties of glass/epoxy and carbon/epoxy							
Material constants	E <sub>1</sub> (GPa)	E <sub>2</sub> (GPa)	G <sub>1</sub> 2 (GPa)	$v_12$	v <sub>2</sub> 3	$\rho(\text{gcm}^{-3})$		
Glass/epoxv	38.49	9.367	3.414	0.2912	0.5071	2.66		
Carbon/epoxy	142.17	9.255	4.795	0.3340	0.4862	1.90		

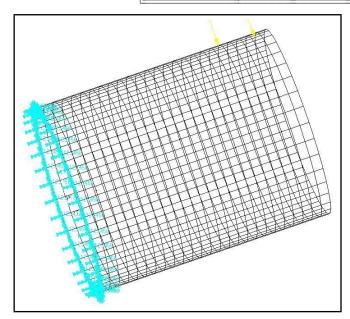


Figure 14. The finite element model of composite laminated cylindrical shell.

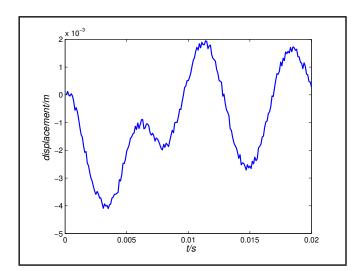


Figure 15. The corresponding radial displacement response.

mal deviation and average deviation in identifying the first impact force are 25.65%, 8.38%, respectively. The average deviation and maximal deviation of MCG in identifying the first impact force are 20.20%, 7.67%, respectively. Additionally, this table also shows that Landweber's and FRCG's average and maximal deviation in the recognition of the second impact force are both larger than MCG's. At the same time, the iterative numbers of Landweber and FRCG are 160 and 108, respectively, while the iterative number of the proposed method is 90. Therefore, these results illustrate that the newly developed MCG method is stable and effective in reconstructing the impact loads.

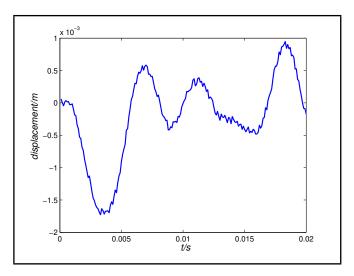
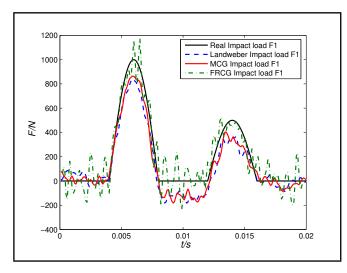
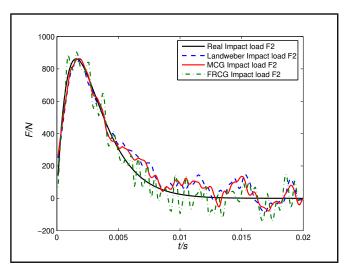


Figure 16. The corresponding radial displacement response.



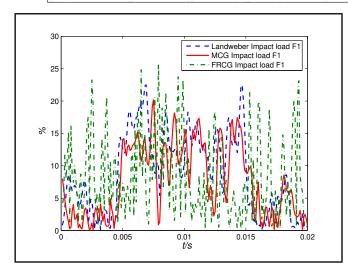
**Figure 17.** The identified first impact load  $F_1$  at noise level 10%.



**Figure 18.** The identified second impact load  $F_2$  at noise level 10%.

**Table 6.** The identified impact force at five time points at noise level 10%.

			Landweber method		LMCG method		LFRCG method	
	Timepoint	Realforce	Identifiedforce	Error (%)	Identifiedforce	Error (%)	Identified force	Error (%)
Impact load F1	0.0005	0	76.20	7.62	22.71	2.27	-109.25	10.93
Impact load F2	0.0004	464.49	408.79	6.46	435.44	3.37	486.7	2.58
Impact load F1	0.006	1000	833.06	16.69	863.05	13.70	1156.8	15.68
Impact load F2	0.0016	861.97	836.04	3.01	857.81	0.48	904.58	4.94
Impact load F1	0.0122	78.22	3.16	7.51	11.88	6.63	66.24	1.20
Impact load F2	0.0026	738.58	701.87	4.26	746.47	0.92	839.36	11.69
Impact load F1	0.014	500	361.71	13.83	352.51	14.75	438.25	6.18
Impact load F2	0.0102	22.37	75.13	6.12	79.55	6.63	-26.59	5.68
Impact load F1	0.0169	0	31.56	3.16	6.26	0.63	187.36	18.74
Impact load F2 0.0126 5.95		5.95	2.33	0.42	59.35	6.20	-70.60	8.88
Error (%)			Maximum	Average	Maximum	Average	Maximum	Average
Impact load F1			22.52	8.84	20.20	7.67	25.65	8.38
Impact load F2			17.09	6.55	15.56	5.82	17.31	6.19
Iterative steps			160		90		108	



**Figure 19.** The relative deviations for the identified first impact load  $F_1$  at noise level 10%.

#### 5. CONCLUSION

Dynamic load identification is an important problem in practical structural engineering. The solution to this problem cannot be directly dealt with by the traditional inverse matrix method. In this paper, we have proposed a new fast convergent conjugate gradient algorithm for the dynamic force reconstruction of an airfoil structure and a composite laminated cylindrical shell. Numerical performances have shown that the proposed algorithm is a powerful method for dynamic load identification.

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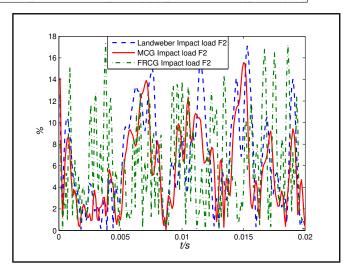


Figure 20. The relative deviations for the identified second impact load  $F_2$  at noise level 10%.

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